

ON ABSENCE OF BOUND STATES FOR WEAKLY ATTRACTIVE δ' -INTERACTIONS SUPPORTED ON NON-CLOSED CURVES IN \mathbb{R}^2

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ABSTRACT. Let $\Lambda \subset \mathbb{R}^2$ be a non-closed piecewise- C^1 curve, which is either bounded with two free endpoints or unbounded with one free endpoint. Let $u_{\pm}|_{\Lambda} \in L^2(\Lambda)$ be the traces of a function u in the Sobolev space $H^1(\mathbb{R}^2 \setminus \Lambda)$ onto two faces of Λ . We prove that for a wide class of shapes of Λ the Schrödinger operator H_{ω}^{Λ} with δ' -interaction supported on Λ of strength $\omega \in L^{\infty}(\Lambda; \mathbb{R})$ associated with the quadratic form

$$H^1(\mathbb{R}^2 \setminus \Lambda) \ni u \mapsto \int_{\mathbb{R}^2} |\nabla u|^2 dx - \int_{\Lambda} \omega |u_+|_{\Lambda} - u_-|_{\Lambda}|^2 ds$$

has no negative spectrum provided that ω is pointwise majorized by a strictly positive function explicitly expressed in terms of Λ . If, additionally, the domain $\mathbb{R}^2 \setminus \Lambda$ is quasi-conical, we show that $\sigma(H_{\omega}^{\Lambda}) = [0, +\infty)$. For a bounded curve Λ in our class and non-varying interaction strength $\omega \in \mathbb{R}$ we derive existence of a constant $\omega_* > 0$ such that $\sigma(H_{\omega}^{\Lambda}) = [0, +\infty)$ for all $\omega \in (-\infty, \omega_*]$; informally speaking, bound states are absent in the weak coupling regime.

1. INTRODUCTION

In this paper we study the self-adjoint operator corresponding to the formal differential expression

$$-\Delta - \omega \delta'(\cdot - \Lambda), \quad \text{on } \mathbb{R}^2,$$

with the δ' -interaction supported on a non-closed piecewise- C^1 curve $\Lambda \subset \mathbb{R}^2$, which is either bounded with two free endpoints or unbounded with one free endpoint, here $\omega \in L^{\infty}(\Lambda; \mathbb{R})$ is called the strength of the interaction. More precisely, for any function u in the Sobolev space $H^1(\mathbb{R}^2 \setminus \Lambda)$ its traces $u_{\pm}|_{\Lambda}$ onto two faces of Λ turn out to be well-defined as functions in $L^2(\Lambda)$, and employing the shorthand notation $[u]_{\Lambda} := u_+|_{\Lambda} - u_-|_{\Lambda}$ we introduce the following symmetric sesquilinear form

$$(1.1) \quad \begin{aligned} \mathfrak{a}_{\omega}^{\Lambda}[u, v] &:= (\nabla u, \nabla v)_{L^2(\mathbb{R}^2; \mathbb{C}^2)} - (\omega[u]_{\Lambda}, [v]_{\Lambda})_{L^2(\Lambda)}, \\ \text{dom } \mathfrak{a}_{\omega}^{\Lambda} &:= H^1(\mathbb{R}^2 \setminus \Lambda), \end{aligned}$$

which is closed, densely defined, and semibounded in the Hilbert space $L^2(\mathbb{R}^2)$; see Proposition 3.1. Let H_{ω}^{Λ} be defined as the unique self-adjoint operator representing the form $\mathfrak{a}_{\omega}^{\Lambda}$ in the usual manner. We regard H_{ω}^{Λ} as the *Schrödinger operator with δ' -interaction of strength ω supported on Λ* .

The aim of this paper is to demonstrate a peculiar spectral property of H_{ω}^{Λ} . Namely, we show absence of negative spectrum for H_{ω}^{Λ} under not too restrictive assumptions on the shape of Λ and assuming that the strength ω is pointwise majorized by a strictly positive function explicitly expressed in terms of the shape of Λ . The important point to note here is that the discovered phenomenon is non-emergent for δ' -interactions supported on loops in \mathbb{R}^2 ; cf. [3, Thm. 4.4].

Key words and phrases. Schrödinger-type operators, δ' -interactions, non-closed curves, negative spectrum, min-max principle, linear fractional transformations.

The basic geometric ingredient in our paper is the concept of *monotone curves*. A non-closed piecewise- C^1 curve $\Lambda \subset \mathbb{R}^2$ is monotone if it can be parametrized via the piecewise- C^1 mapping $\varphi: (0, R) \rightarrow \mathbb{R}$, $R \in (0, +\infty]$, as

$$(1.2) \quad \Lambda = \{x_0 + (r \cos \varphi(r), r \sin \varphi(r)) \in \mathbb{R}^2 : r \in (0, R)\};$$

here $x_0 \in \mathbb{R}^2$ is fixed. For example, a circular arc subtending an angle $\theta \leq \pi$ is monotone, whereas a circular arc subtending an angle $\theta > \pi$ is not.

In the next theorem, which is the first main result of our paper, we provide a condition on ω ensuring absence of negative spectrum for the operator H_ω^Λ with Λ being monotone. The statement of Theorem A below is contained in Theorem 4.2, in Subsection 4.2.

Theorem A. *Let a monotone piecewise- C^1 curve $\Lambda \subset \mathbb{R}^2$ be parametrized as in (1.2) via $\varphi: (0, R) \rightarrow \mathbb{R}$, $R \in (0, +\infty]$. Then the spectrum of H_ω^Λ satisfies*

$$\sigma(H_\omega^\Lambda) \subseteq [0, +\infty) \quad \text{if} \quad \omega(r) \leq \frac{1}{2\pi r \sqrt{1 + (r\varphi'(r))^2}}, \quad \text{for } r \in (0, R).$$

If ω is majorized as above and, additionally, the domain $\mathbb{R}^2 \setminus \Lambda$ is quasi-conical, then $\sigma(H_\omega^\Lambda) = [0, +\infty)$.

Roughly speaking, a domain $\Omega \subset \mathbb{R}^2$ is quasi-conical if it contains a disc of arbitrary large radius; see Subsection 3.2 for details. In Proposition 4.7 we demonstrate that, in general, the operator H_ω^Λ may have negative spectrum if the δ' -interaction is “sufficiently strong”.

Operators H_ω^Λ with non-varying strengths $\omega \in \mathbb{R}$ are of special interest. One can derive from Theorem A that for a bounded monotone Λ one can find a constant $\omega_* > 0$ such that

$$(1.3) \quad \sigma(H_\omega^\Lambda) = [0, +\infty) \quad \text{for all } \omega \in (-\infty, \omega_*];$$

in other words, there are no bound states in the weak coupling regime. Computation of the largest constant $\omega_* > 0$ such that (1.3) still holds presents a more delicate problem, which will be considered elsewhere.

In the formulation of the second main result of the paper we use the notion of a *linear fractional transformation (LFT)*. The complex plane \mathbb{C} can be extended up to the *Riemann sphere* $\widehat{\mathbb{C}} := \mathbb{C} \cup \{\infty\}$ with a suitable topology and for $a, b, c, d \in \mathbb{C}$ such that $ad - bc \neq 0$ one defines the LFT as

$$M: \widehat{\mathbb{C}} \rightarrow \widehat{\mathbb{C}}, \quad M(z) := \frac{az + b}{cz + d},$$

with the exception of the points $z = \infty$ and $z = -d/c$ if $c \neq 0$, which have to be treated separately; see Subsection 2.3. The next theorem generalizes Theorem A to the case of curves, which are images of monotone curves under LFTs; the statement of this theorem is contained in Theorem 4.12, in Subsection 4.3. Here, we confine ourselves to non-varying interaction strengths only.

Theorem B. *Let $\Lambda \subset \mathbb{R}^2$ be a bounded piecewise- C^1 curve. Suppose that there exists an LFT $M: \widehat{\mathbb{C}} \rightarrow \widehat{\mathbb{C}}$ such that $M(\infty), M^{-1}(\infty) \notin \Lambda$ and that $M^{-1}(\Lambda)$ is monotone. Then there exists a constant $\omega_* > 0$ such that*

$$\sigma(H_\omega^\Lambda) = [0, +\infty) \quad \text{for all } \omega \in (-\infty, \omega_*].$$

In the main body of the paper also an explicit formula for ω_* in the above theorem is provided. Using Theorem B we can treat, *e.g.*, any circular arc, since it can be mapped via a suitable LFT to a subinterval of the straight line in \mathbb{R}^2 ; see Example 4.13. One may even conjecture that for any bounded Λ there exists an $\omega_* > 0$ such that $\sigma(H_\omega^\Lambda) = [0, +\infty)$ for all $\omega \in (-\infty, \omega_*]$.

Our proofs rely on the min-max principle applied to the form $\mathfrak{a}_\omega^\Lambda$ in (1.1) on a suitable core. A further important ingredient in our analysis is careful investigation of a model one-dimensional problem with a point δ' -interaction on the loop.

The results of this paper contribute to a prominent topic in spectral theory: existence/non-existence of weakly coupled bound states for Schrödinger-type operators. Absence of bound states in the weak coupling regime holds for Schrödinger operators with regular potentials in space dimensions $d \geq 3$ (but not for $d = 1, 2$!); see [38]. Also such an effect occurs for δ -interactions supported on arbitrary compact hypersurfaces in \mathbb{R}^3 (see [13]) and for δ -interactions on compact non-closed curves in \mathbb{R}^3 (see [17]). However, for δ -interactions in \mathbb{R}^2 supported on arbitrary compact curves such an effect is non-existent [19, 26].

Schrödinger operators with δ' -interactions supported on hypersurfaces are attractive from physical point of view, because they exhibit rather unusual scattering properties; cf. [1, Chap. 1.4]. These operators are also physically relevant in photonic crystals theory [20]. As a mathematical abstraction they were perhaps first studied in [2, 37], where interactions were supported on spheres. A rigorous definition of such operators with interactions supported on general hypersurfaces has been posed in [12, §7.2] as an open question. Such Hamiltonians with interactions supported on closed hypersurfaces without free boundaries have been rigorously defined in [5] using two approaches: via the theory of self-adjoint extensions of symmetric operators and by means of form methods. Spectral properties of them were investigated in several subsequent works [3, 4, 14, 15, 16, 24, 32]. In the recent preprint [33] Schrödinger operators with δ' -interactions supported on non-closed curves and surfaces are defined via the theory of self-adjoint extensions and their scattering properties are discussed.

Let us briefly outline the structure of the paper. Section 2 presents some preliminaries: Sobolev spaces, geometry of curves, linear fractional transformations, and a model one-dimensional spectral problem. Section 3 provides a rigorous definition of the operator H_ω^Λ and a characterisation of its essential spectrum. Section 4 contains proofs of our main results, formulated in Theorems A and B, as well as some related results and examples. In Section 5 final remarks are given and two open questions are posed. A couple of standard proofs of identities related to LFTs are outsourced to Appendix A.

Notations. By $D_R(\mathbf{x}) := \{\mathbf{x} \in \mathbb{R}^2 : |\mathbf{x} - \mathbf{x}_0| < R\}$ we denote the open disc of the radius $R > 0$ with the center $\mathbf{x}_0 \in \mathbb{R}^2$. If such a disc is centered at the origin, we use the shorthand notation $D_R := D_R(0)$. By definition we set $D_\infty := \mathbb{R}^2$. For a self-adjoint operator T we denote by $\sigma_{\text{ess}}(T)$, $\sigma_d(T)$, and $\sigma(T)$ its essential, discrete, and full spectra, respectively. For an open set $\Omega \subset \mathbb{R}^2$ the space of smooth compactly supported functions and the first order order Sobolev space are denoted by $\mathcal{D}(\Omega)$ and by $H^1(\Omega)$, respectively.

Acknowledgements. The authors are indebted to Jussi Behrndt and Pavel Exner for their active interest in the preparation of this paper and for many stimulating conversations. David Krejčířík is acknowledged for a comment, which led to an improvement in the definition of the operator. MJ was supported by the Grant Agency of the Czech Technical University in Prague, grant No. SGS13/217/OHK4/3T/14. VL was supported by the Austrian Science Fund (FWF) under the project P 25162-N26. Both the authors acknowledge the financial support by the Austria-Czech Republic co-operation grant CZ01/2013 and by the Czech Science Foundation (GAČR) under the project 14-06818S.

2. PRELIMINARIES

This section contains some preliminary material that will be used in the main part of this paper. In Subsection 2.1 we provide basic facts on the Sobolev space H^1 , in particular, we define the Sobolev space $H^1(\mathbb{R}^2 \setminus \Lambda)$ for a non-closed Lipschitz curve Λ . In Subsection 2.2 we introduce the notions of a piecewise- C^1 curve and of a monotone curve. The concept of the linear fractional transformation is discussed in Subsection 2.3. A model spectral problem for one-dimensional Schrödinger operator with one-center δ' -interaction on a loop is considered in Subsection 2.4 and a sufficient condition for absence of negative eigenvalues in this spectral problem is established.

2.1. Sobolev spaces. Let $\Omega \subset \mathbb{R}^2$ be a simply connected *Lipschitz domain* from the class described in [39, Ch. VI]. This class of Lipschitz domains includes (as a subclass) Lipschitz domains with compact boundaries as in [35, Ch. 3], *hypographs* of uniformly Lipschitz functions, and some other domains with non-compact boundaries. In what follows the Hilbert spaces $L^2(\Omega)$, $L^2(\Omega; \mathbb{C}^2)$, $L^2(\partial\Omega)$, and $H^1(\Omega)$ are defined as usual; see *e.g.* [35, Ch. 3] and [34]. For the sake of brevity we denote the scalar products in both $L^2(\Omega)$ and $L^2(\Omega; \mathbb{C}^2)$ by $(\cdot, \cdot)_\Omega$ without any danger of confusion. The scalar product in $L^2(\partial\Omega)$ is abbreviated by $(\cdot, \cdot)_{\partial\Omega}$. The space of functions on $\bar{\Omega}$ smooth up to the boundary $\partial\Omega$ is defined as

$$\mathcal{D}(\bar{\Omega}) := \{u|_\Omega : u \in \mathcal{D}(\mathbb{R}^2)\}.$$

By [35, Thm. 3.29], see also [34, 39], the space $\mathcal{D}(\bar{\Omega})$ is dense in both $L^2(\Omega)$ and $H^1(\Omega)$. The natural restriction mapping $\mathcal{D}(\bar{\Omega}) \ni u \mapsto u|_{\partial\Omega} \in L^2(\partial\Omega)$ can be extended by continuity up to the whole space $H^1(\Omega)$; see *e.g.* [35, Thm. 3.37] and [34]. The corresponding extension by continuity $H^1(\Omega) \ni u \mapsto u|_{\partial\Omega} \in L^2(\partial\Omega)$ is called *the trace mapping*. The statement of the first lemma in this subsection appears in several monographs and papers in various forms; see *e.g.* [3, Lem. 2.6] and [21, Lem. 2.5] for two different proofs of this statement.

Lemma 2.1. *Let $\Omega \subset \mathbb{R}^2$ be a Lipschitz domain. Then for any $\varepsilon > 0$ there exists a constant $C(\varepsilon) > 0$ such that*

$$\|u|_{\partial\Omega}\|_{\partial\Omega}^2 \leq \varepsilon \|\nabla u\|_\Omega^2 + C(\varepsilon) \|u\|_\Omega^2$$

holds for all $u \in H^1(\Omega)$.

The following hypothesis will be used throughout the paper.

Hypothesis 2.1. *Let $\Omega_+ \subset \mathbb{R}^2$ be a simply connected Lipschitz domain from the above class, whose complement $\Omega_- := \mathbb{R}^2 \setminus \bar{\Omega}_+$ is a Lipschitz domain from the same class. Set $\Sigma := \partial\Omega_+ = \partial\Omega_-$ and suppose that $\Lambda \subset \Sigma$ is a connected subarc of Σ , which is not necessarily bounded if Σ is unbounded.*

Obviously, the orthogonal sum $H^1(\Omega_+) \oplus H^1(\Omega_-)$ is a Hilbert space with respect to the scalar product

$$(u_+ \oplus u_-, v_+ \oplus v_-)_1 := (u_+, v_+)_{H^1(\Omega_+)} + (u_-, v_-)_{H^1(\Omega_-)}, \quad u_\pm, v_\pm \in H^1(\Omega_\pm).$$

The norm associated to this scalar product is denoted by $\|\cdot\|_1$. Let us define the jump of the trace as

$$[u]_\Sigma := u_+|_\Sigma - u_-|_\Sigma, \quad u = u_+ \oplus u_- \in H^1(\Omega_+) \oplus H^1(\Omega_-).$$

The Hilbert space $L^2(\Sigma)$ can be decomposed into the orthogonal sum

$$L^2(\Sigma) = L^2(\Lambda) \oplus L^2(\Sigma \setminus \Lambda).$$

The scalar products in $L^2(\Lambda)$ and $L^2(\Sigma \setminus \Lambda)$ will further be denoted by $(\cdot, \cdot)_\Lambda$ and $(\cdot, \cdot)_{\Sigma \setminus \Lambda}$. Clearly enough, the restrictions of $u_\pm|_\Sigma$ for a $u_\pm \in H^1(\Omega_\pm)$ to the arcs $\Sigma \setminus \Lambda$ and Λ satisfy $u_\pm|_{\Sigma \setminus \Lambda} \in L^2(\Sigma \setminus \Lambda)$ and $u_\pm|_\Lambda \in L^2(\Lambda)$. Let us also introduce the notations

$$[u]_\bullet := u_+|_\bullet - u_-|_\bullet, \quad \bullet = \Lambda, \Sigma \setminus \Lambda, \quad u = u_+ \oplus u_- \in H^1(\Omega_+) \oplus H^1(\Omega_-).$$

The linear space

$$(2.1) \quad F_\Lambda := \{u \in \mathcal{D}(\overline{\Omega_+}) \oplus \mathcal{D}(\overline{\Omega_-}) : [u]_{\Sigma \setminus \Lambda} = 0\}$$

is a subspace of the Hilbert space $H^1(\Omega_+) \oplus H^1(\Omega_-)$, and its closure in $H^1(\Omega_+) \oplus H^1(\Omega_-)$

$$(2.2) \quad H^1(\mathbb{R}^2 \setminus \Lambda) := \overline{F_\Lambda}^{\|\cdot\|_1}$$

is itself a Hilbert space with respect to the same scalar product $(\cdot, \cdot)_1$.

Remark 2.2. The above construction of the space $H^1(\mathbb{R}^2 \setminus \Lambda)$ can easily be translated to the higher space dimensions, in which case Λ will be a hypersurface with free boundary (open hypersurface).

Remark 2.3. The space $H^1(\mathbb{R}^2 \setminus \Lambda)$ can also be defined in an alternative way. The set $\mathbb{R}^2 \setminus \Lambda$ is an open subset of \mathbb{R}^2 . Hence, one can define for any $u \in L^2(\mathbb{R}^2)$ its weak partial derivatives $\partial_1 u$ and $\partial_2 u$ by means of the test functions in $\mathcal{D}(\mathbb{R}^2 \setminus \Lambda)$; see *e.g.* [35, Ch. 3]. Then the space $H^1(\mathbb{R}^2 \setminus \Lambda)$ is given by

$$H^1(\mathbb{R}^2 \setminus \Lambda) = \{u \in L^2(\mathbb{R}^2) : \partial_1 u, \partial_2 u \in L^2(\mathbb{R}^2)\}.$$

We are not aiming to provide here an argumentation that this new definition gives rise to the same space as in (2.2). It is only important here that the equivalence of these definitions automatically implies that the space $H^1(\mathbb{R}^2 \setminus \Lambda)$ is independent of the continuation of the arc Λ up to Σ . Another way of verifying the independence of the space $H^1(\mathbb{R}^2 \setminus \Lambda)$ from a continuation of Λ can be found in [9].

Next proposition collects some useful properties of the traces of functions in $H^1(\mathbb{R}^2 \setminus \Lambda)$ onto $\Sigma \setminus \Lambda$ and onto Λ .

Proposition 2.4. *Let the curves $\Sigma, \Lambda \subset \mathbb{R}^2$, and the domains $\Omega_\pm \subset \mathbb{R}^2$ be as in Hypothesis 2.1. Let the Hilbert space $(H^1(\mathbb{R}^2 \setminus \Lambda), (\cdot, \cdot)_1)$ be as in (2.2). Then the following statements hold.*

- (i) $[u]_{\Sigma \setminus \Lambda} = 0$ for all $u \in H^1(\mathbb{R}^2 \setminus \Lambda)$.
- (ii) For any $\varepsilon > 0$ there exists a constant $C(\varepsilon) > 0$ such that

$$\|[u]_\Lambda\|_\Lambda^2 \leq \varepsilon \|\nabla u\|_{\mathbb{R}^2}^2 + C(\varepsilon) \|u\|_{\mathbb{R}^2}^2$$

for all $u \in H^1(\mathbb{R}^2 \setminus \Lambda)$.

Proof. (i) It can be easily checked that the continuity of the trace mappings

$$H^1(\Omega_\pm) \ni u_\pm \mapsto u_\pm|_\Sigma \in L^2(\Sigma)$$

implies that the mapping

$$H^1(\Omega_+) \oplus H^1(\Omega_-) \ni u \mapsto [u]_{\Sigma \setminus \Lambda} \in L^2(\Sigma \setminus \Lambda)$$

is well-defined and continuous. Note that for any $u \in H^1(\mathbb{R}^2 \setminus \Lambda)$ there exists an approximating sequence $(u_n)_n \subset F_\Lambda$ (*cf.* (2.2)) such that $\|u_n - u\|_1 \rightarrow 0$ as $n \rightarrow \infty$. Hence, we obtain

$$[u]_{\Sigma \setminus \Lambda} = \lim_{n \rightarrow \infty} [u_n]_{\Sigma \setminus \Lambda} = 0.$$

(ii) By Lemma 2.1 for any $\varepsilon > 0$ there exist constants $C_{\pm}(\varepsilon) > 0$ such that

$$(2.3) \quad \|u_{\pm}|_{\Sigma}\|_{\Sigma}^2 \leq (\varepsilon/2)\|\nabla u_{\pm}\|_{\Omega_{\pm}}^2 + C_{\pm}(\varepsilon)\|u_{\pm}\|_{\Omega_{\pm}}^2$$

for all $u \in H^1(\Omega_+) \oplus H^1(\Omega_-)$. Set then $C(\varepsilon) := \max\{2C_+(\varepsilon), 2C_-(\varepsilon)\}$. Using the result of item (i) and the bounds (2.3) we obtain that for any $\varepsilon > 0$ and any $u = u_+ \oplus u_- \in H^1(\mathbb{R}^2 \setminus \Lambda) \subset H^1(\Omega_+) \oplus H^1(\Omega_-)$ holds

$$\begin{aligned} \|[u]_{\Lambda}\|_{\Lambda}^2 &= \|[u]_{\Sigma}\|_{\Sigma}^2 \leq 2\|u_+|_{\Sigma}\|_{\Sigma}^2 + 2\|u_-|_{\Sigma}\|_{\Sigma}^2 \\ &\leq \varepsilon\|\nabla u_+\|_{\Omega_+}^2 + \varepsilon\|\nabla u_-\|_{\Omega_-}^2 + 2C_+(\varepsilon)\|u_+\|_{\Omega_+}^2 + 2C_-(\varepsilon)\|u_-\|_{\Omega_-}^2 \\ &\leq \varepsilon\|\nabla u\|_{\mathbb{R}^2}^2 + C(\varepsilon)\|u\|_{\mathbb{R}^2}^2. \end{aligned}$$

□

Remark 2.5. For $\omega_1, \omega_2 \in L^\infty(\Lambda; \mathbb{R})$ by writing $\omega_1 \leq \omega_2$ we will always implicitly mean that $\omega_2 - \omega_1 \geq 0$ almost everywhere.

2.2. On curves in \mathbb{R}^2 . We begin this subsection by defining the notion of a piecewise- C^1 curve. It should be emphasized that, especially for unbounded curves, definition of a piecewise- C^1 curve is non-unique in the mathematical literature.

Definition 2.6. A non-closed curve $\Lambda \subset \mathbb{R}^2$ satisfying Hypothesis 2.1 is called piecewise- C^1 if it can be parametrized via a piecewise- C^1 mapping

$$(2.4) \quad \lambda: I \rightarrow \mathbb{R}^2, \quad \lambda(s) := (\lambda_1(s), \lambda_2(s)), \quad I := (0, L), \quad L \in (0, +\infty],$$

such that $\lambda(I) = \Lambda$ and λ is injective. If, moreover, $|\lambda'(s)| = 1$ for almost all $s \in I$, then such a parametrization is called natural and L is then called the length of Λ .

We require in the above definition, that the curve Λ satisfies Hypothesis 2.1, to avoid increasing oscillations at infinity for unbounded curves.

Further, we proceed to define a (non-standard) concept of a *monotone curve*. The authors have not succeeded to find a common name for this concept in the literature on geometry.

Definition 2.7. A piecewise- C^1 curve $\Lambda \subset \mathbb{R}^2$ is called monotone if it can be parametrized via a piecewise- C^1 mapping $\varphi: (0, R) \rightarrow \mathbb{R}$ with $R \in (0, +\infty]$ such that

$$\Lambda = \{x_0 + (r \cos \varphi(r), r \sin \varphi(r)) \in \mathbb{R}^2 : r \in (0, R)\}$$

with some fixed $x_0 \in \mathbb{R}^2$.

Informally speaking, a curve Λ is monotone if the distance (measured in \mathbb{R}^2) from one of its endpoints is always increasing when travelling along Λ from this endpoint towards another endpoint or towards infinity.

Remark 2.8. For a curve $\Lambda \subset \mathbb{R}^2$ as in Definition 2.7 any function $\omega \in L^\infty(\Lambda)$ can be viewed as a function of the argument $r \in (0, R)$.

2.3. Linear fractional transformations. For later purposes we introduce linear fractional transformations (LFT) and state several useful properties of them. To work with LFT it is more convenient to deal with the *extended complex plane (Riemann sphere)* $\widehat{\mathbb{C}} := \mathbb{C} \cup \{\infty\}$ rather than the usual complex plane. The complex plane itself as a subset of $\widehat{\mathbb{C}}$ can be naturally identified with the Euclidean plane \mathbb{R}^2 and occasionally we will use this identification.

For the purpose of convenience the extended complex plane $\widehat{\mathbb{C}}$ is endowed with a suitable topology: a sequence $(z_n)_n \in \widehat{\mathbb{C}}$ converges to $z \in \widehat{\mathbb{C}}$ if one of the following conditions holds:

- (i) $z = \infty$ and there exists $N \in \mathbb{N}$ such that $z_n = \infty$ for all $n \geq N$;
- (ii) $z = \infty$ and any infinite subsequence $(z_{n_k})_k \subset \mathbb{C}$ of $(z_n)_n$ satisfies $\lim_{k \rightarrow \infty} |z_{n_k}| = \infty$;
- (iii) $z \in \mathbb{C}$, there exists $N \in \mathbb{N}$ such that $z_n \neq \infty$ for all $n \geq N$, and $\lim_{n \rightarrow \infty} z_n = z$ in the sense of convergence in \mathbb{C} .

This definition of topology can also be easily reformulated in terms of open sets. The above topology on $\hat{\mathbb{C}}$ is equivalent to the topology of \mathbb{S}^2 (unit sphere in \mathbb{R}^3). A natural homeomorphism between $\hat{\mathbb{C}}$ and \mathbb{S}^2 is called *stereographic projection*; see e.g. [29, §6.2.3].

For $a, b, c, d \in \mathbb{C}$ such that $ad - bc \neq 0$ the mapping $M: \hat{\mathbb{C}} \rightarrow \hat{\mathbb{C}}$ is an LFT if one of the two conditions holds:

- (i) $c = 0$, $d \neq 0$, $M(\infty) := \infty$, and $M(z) := (a/d)z + (b/d)$ for $z \in \mathbb{C}$.
- (ii) $c \neq 0$, $M(\infty) := a/c$, $M(-d/c) := \infty$, and $M(z) := \frac{az+b}{cz+d}$ for $z \in \mathbb{C}$, $z \neq -d/c$.

The following statement can be found in [29, §6.2.3].

Proposition 2.9. *Any LFT $M: \hat{\mathbb{C}} \rightarrow \hat{\mathbb{C}}$ is a homeomorphism with respect to the above topology on $\hat{\mathbb{C}}$ and its inverse M^{-1} is also an LFT. The composition $M_1 \circ M_2$ of two LFTs M_1, M_2 is an LFT as well.*

It is convenient to introduce $M_1(x, y) := \operatorname{Re} M(x + iy)$ and $M_2(x, y) := \operatorname{Im} M(x + iy)$. Then *Cauchy-Riemann equations*

$$(2.5) \quad \partial_x M_1 = \partial_y M_2, \quad \partial_x M_2 = -\partial_y M_1,$$

hold pointwise in \mathbb{R}^2 except the point $M^{-1}(\infty)$. In view of these equations the Jacobian J_M of the mapping M can be computed (again except the point $M^{-1}(\infty)$) by the formulae

$$(2.6) \quad J_M = (\partial_x M_1)^2 + (\partial_y M_1)^2 = (\partial_x M_2)^2 + (\partial_y M_2)^2;$$

also the following relation turns out to be useful

$$(2.7) \quad \langle \nabla M_1, \nabla M_2 \rangle = \partial_x M_1 \partial_x M_2 + \partial_y M_1 \partial_y M_2 = 0;$$

i.e. the vectors ∇M_1 and ∇M_2 are orthogonal to each other.

Next auxiliary lemma is of purely technical nature and is proven for convenience of the reader.

Lemma 2.10. *Let M be an LFT with the Jacobian J_M . Then for any $x \in \mathbb{R}^2$, $x \neq M^{-1}(\infty)$, and any function $u: \mathbb{R}^2 \simeq \mathbb{C} \rightarrow \mathbb{C}$ differentiable at the point $M(x)$*

$$|(\nabla v)(x)|^2 = |(\nabla u)(M(x))|^2 J_M(x)$$

holds with $v = u \circ M$.

Proof. Using relations (2.6), (2.7), and the chain rule for differentiation we obtain

$$\begin{aligned} |\nabla v|^2 &= |(u'_x \circ M) \partial_x M_1 + (u'_y \circ M) \partial_x M_2|^2 + |(u'_x \circ M) \partial_y M_1 + (u'_y \circ M) \partial_y M_2|^2 \\ &= (|u'_x \circ M|^2 + |u'_y \circ M|^2) J_M + 2 \operatorname{Re} \left[((u'_x u'_y) \circ M) \cdot \langle \nabla M_1, \nabla M_2 \rangle \right] \\ &= (|u'_x \circ M|^2 + |u'_y \circ M|^2) J_M = |(\nabla u) \circ M|^2 J_M. \end{aligned}$$

The claim is thus shown. \square

2.4. Point δ' -interaction on a loop. In this subsection we introduce an auxiliary self-adjoint Schrödinger operator $\mathsf{T}_{d,\omega}$ acting in the Hilbert space $(L^2(I), (\cdot, \cdot)_I)$ with $I := (0, d)$ and corresponding to a point δ' -interaction on the one-dimensional loop of length $d > 0$. Employing the following shorthand notation

$$[\psi]_{\partial I} := \psi(d-) - \psi(0+), \quad \psi \in H^2(I),$$

we define

$$(2.8) \quad \mathsf{T}_{d,\omega}\psi := -\psi'', \quad \text{dom } \mathsf{T}_{d,\omega} := \{\psi \in H^2(I) : \psi'(0+) = \psi'(d-) = \omega[\psi]_{\partial I}\},$$

where $\omega \in \mathbb{R}$; see [1, 7, 10, 18, 23, 28] for the investigations of more general operators of this type. Note that $\omega = 0$ corresponds to Neumann boundary conditions at the endpoints. Next proposition states a spectral property of $\mathsf{T}_{d,\omega}$, which is useful for our purposes.

Proposition 2.11. *The self-adjoint operator $\mathsf{T}_{d,\omega}$ in the Hilbert space $L^2(I)$, defined in (2.8), is non-negative if $d\omega \leq 1$.*

Proof. We prove this proposition via construction of an explicit condition for the negative spectrum of $\mathsf{T}_{d,\omega}$ and its analysis. Obviously, the spectrum of $\mathsf{T}_{d,\omega}$ is discrete (due to the compact embedding of $H^2(I)$ into $L^2(I)$). An eigenfunction of $\mathsf{T}_{d,\omega}$, which corresponds to a negative eigenvalue $\lambda = -\kappa^2 < 0$ ($\kappa > 0$) is characterized by the following two conditions:

$$(2.9a) \quad -\psi''(x) = -\kappa^2\psi(x);$$

$$(2.9b) \quad \psi'(0+) = \psi'(d-) = \omega[\psi]_{\partial I}.$$

The condition (2.9a) is satisfied by a function, which can be represented in the form

$$\psi(x) = A \exp(\kappa x) + B \exp(-\kappa x), \quad x \in (0, d),$$

with some constants $A, B \in \mathbb{C}$. Simple computations yield

$$\begin{aligned} \psi(0+) &= A + B, & \psi(d-) &= A \exp(\kappa d) + B \exp(-\kappa d), \\ \psi'(0+) &= \kappa A - \kappa B, & \psi'(d-) &= \kappa A \exp(\kappa d) - \kappa B \exp(-\kappa d). \end{aligned}$$

The above identities and the condition (2.9b) together imply

$$(2.10a) \quad A = \frac{1 - \exp(-\kappa d)}{1 - \exp(\kappa d)} B;$$

$$(2.10b) \quad \kappa A - \kappa B = \omega \left(A(\exp(\kappa d) - 1) + B(\exp(-\kappa d) - 1) \right).$$

Substituting the formula (2.10a) into (2.10b), we arrive at

$$\kappa B \left(\frac{1 - \exp(-\kappa d)}{1 - \exp(\kappa d)} - 1 \right) = \omega \left(-B(1 - \exp(-\kappa d)) + B(\exp(-\kappa d) - 1) \right),$$

that is equivalent to

$$\exp(-\kappa d) - \exp(\kappa d) = \frac{2\omega}{\kappa} \left(1 - \exp(-\kappa d) \right) \left(1 - \exp(\kappa d) \right).$$

Making several steps further in the computations, we obtain

$$\begin{aligned} 1 &= \frac{2\omega}{\kappa} \frac{(1 - \exp(-\kappa d))(1 - \exp(\kappa d))}{\exp(\kappa d)(\exp(-2\kappa d) - 1)} = \frac{2\omega}{\kappa} \frac{\exp(\kappa d) - 1}{\exp(\kappa d)(\exp(-\kappa d) + 1)} \\ &= \frac{2\omega}{\kappa} \frac{1 - \exp(-\kappa d)}{1 + \exp(-\kappa d)}. \end{aligned}$$

Define then the following function

$$\Theta_\omega(\kappa) := \frac{2\omega}{\kappa} \frac{1 - \exp(-\kappa d)}{1 + \exp(-\kappa d)}, \quad \kappa > 0.$$

Hence, the point $\lambda = -\kappa^2$ is a negative eigenvalue of $\mathsf{T}_{d,\omega}$ if and only if $\Theta_\omega(\kappa) = 1$. Let us consider the following auxiliary function

$$f(x) := \frac{1 - e^{-x}}{1 + e^{-x}}, \quad x \geq 0,$$

which is clearly continuously differentiable, and whose derivative is given by

$$f'(x) = \frac{2}{(e^{x/2} + e^{-x/2})^2}, \quad x \geq 0.$$

Hence, using the standard inequality $a + 1/a > 2$, $a \in (0, +\infty)$, $a \neq 1$, we get $f'(x) < 1/2$ for all $x > 0$. Applying the mean value theorem to f , we obtain

$$f(x) = f(0) + f'(\xi)(x - 0) = f'(\xi)x < \frac{x}{2};$$

here $\xi \in (0, x)$. Finally, note that

$$0 \leq \Theta_\omega(\kappa) = \frac{2\omega}{\kappa} f(\kappa d) < d\omega.$$

Thus, for $d\omega \leq 1$ the equation $\Theta_\omega(\kappa) = 1$ has no positive roots and the claim follows. \square

According to *e.g.* [27], the operator $\mathsf{T}_{d,\omega}$ represents the sesquilinear form

$$(2.11) \quad \mathfrak{a}_{d,\omega}[\psi, \varphi] := (\psi', \varphi')_I - \omega[\psi]_{\partial I}[\overline{\varphi}]_{\partial I}, \quad \text{dom } \mathfrak{a}_{d,\omega} := H^1(I),$$

and we can derive the following simple corollary of Proposition 2.11.

Corollary 2.12. *Let the sesquilinear form $\mathfrak{a}_{d,\omega}$ be as in (2.11). If $d\omega \leq 1$, then $\mathfrak{a}_{d,\omega}[\psi] \geq 0$ for all $\psi \in H^1(I)$.*

Remark 2.13. Consider the non-negative symmetric operator

$$S\psi := -\psi'', \quad \text{dom } S := H_0^2(I),$$

in $L^2(I)$. The operator S is known to have deficiency indices $(2, 2)$. One may consider self-adjoint extensions of S in $L^2(I)$. The self-adjoint operator $\mathsf{T}_{d,\omega}$ with $d\omega = 1$ turns to be the *Krein-von Neumann extension* of S (the “smallest” non-negative self-adjoint extension of S); *i.e.* for any other non-negative self-adjoint extension T of S

$$(T + a)^{-1} \leq (\mathsf{T}_{d,\omega} + a)^{-1}$$

holds for all $a > 0$; see *e.g.* [36, Cor. 10.13, Thm. 14.25, Ex. 14.14].

3. DEFINITION OF THE OPERATOR AND ITS ESSENTIAL SPECTRUM

In this section we rigorously define using form methods Schrödinger operators with δ' -interactions supported on non-closed curves as in Hypothesis 2.1 and characterise their essential spectra. In the latter characterisation the notion of a quasi-conical domain plays an essential role.

3.1. Definition of the operator via its sesquilinear form. Schrödinger operators with δ' -interactions supported on closed hypersurfaces were defined and investigated in [2, 3, 5, 14, 15, 16]. The goal of this subsection is to define rigorously Schrödinger operator with δ' -interactions supported on a non-closed curve Λ satisfying Hypothesis 2.1. In the case of a bounded $C^{2,1}$ -smooth curve Λ our definition of the operator agrees with the one in the recent preprint [33], where this Hamiltonian is defined using the theory of self-adjoint extensions of symmetric operators.

Let $\omega \in L^\infty(\Lambda; \mathbb{R})$ and denote by $\|\omega\|_\infty$ its sup-norm. Recall the definition of the sesquilinear form $\mathfrak{a}_\omega^\Lambda$ in (1.1)

$$(3.1) \quad \mathfrak{a}_\omega^\Lambda[u, v] := (\nabla u, \nabla v)_{\mathbb{R}^2} - (\omega[u]_\Lambda, [v]_\Lambda)_\Lambda, \quad \text{dom } \mathfrak{a}_\omega^\Lambda := H^1(\mathbb{R}^2 \setminus \Lambda).$$

If $\omega \equiv 0$, we occasionally write $\mathfrak{a}_\Lambda^\Lambda$ instead of $\mathfrak{a}_\omega^\Lambda$.

Proposition 3.1. *Let $\Lambda \subset \mathbb{R}^2$ be as in Hypothesis 2.1, let $\omega \in L^\infty(\Lambda; \mathbb{R})$, and let the linear space F_Λ be as in (2.1). Then the sesquilinear form $\mathfrak{a}_\omega^\Lambda$ in (3.1) is closed, densely defined, symmetric, and lower-semibounded in the Hilbert space $L^2(\mathbb{R}^2)$. Moreover, $F_\Lambda \subset \text{dom } \mathfrak{a}_\omega^\Lambda$ is a core for this form.*

Proof. Since $\mathfrak{a}_\omega^\Lambda[u, u] \in \mathbb{R}$ for all $u \in \text{dom } \mathfrak{a}_\omega^\Lambda$, the form $\mathfrak{a}_\omega^\Lambda$ is, clearly, symmetric. It is straightforward to see the chain of inclusions $\mathcal{D}(\mathbb{R}^2) \subset F_\Lambda \subset \text{dom } \mathfrak{a}_\omega^\Lambda$. Density of $\text{dom } \mathfrak{a}_\omega^\Lambda$ in $L^2(\mathbb{R}^2)$ follows thus from the density of $\mathcal{D}(\mathbb{R}^2)$ in $L^2(\mathbb{R}^2)$; for the latter see *e.g.* [35, Cor. 3.5].

The norm induced in the conventional way by the form $\mathfrak{a}_\omega^\Lambda$ on its domain $H^1(\mathbb{R}^2 \setminus \Lambda)$ is easily seen to be equivalent to the norm $\|\cdot\|_1$ introduced in Subsection 2.1. Hence, the form $\mathfrak{a}_\omega^\Lambda$ is closed and the space F_Λ , being dense in $H^1(\mathbb{R}^2 \setminus \Lambda)$, is a core for it; *cf.* [36, Dfn. 10.2]. Let us then introduce an auxiliary form

$$\mathfrak{a}'[u, v] := (\omega[u]_\Lambda, [v]_\Lambda)_\Lambda, \quad \text{dom } \mathfrak{a}' := H^1(\mathbb{R}^2 \setminus \Lambda).$$

Using Proposition 2.4 (ii) we get for all $\varepsilon > 0$ the following bound

$$|\mathfrak{a}'[u, u]| \leq \varepsilon \|\omega\|_\infty \mathfrak{a}_\Lambda^\Lambda[u, u] + C(\varepsilon) \|\omega\|_\infty \|u\|_{\mathbb{R}^2}^2$$

with some $C(\varepsilon) > 0$. Choosing $\varepsilon < \frac{1}{\|\omega\|_\infty}$ in the above bound, we obtain that \mathfrak{a}' is relatively bounded with respect to $\mathfrak{a}_\Lambda^\Lambda$ with form bound < 1 . Hence, by [36, Thm. 10.21] (*KLMN theorem*) the form $\mathfrak{a}_\omega^\Lambda = \mathfrak{a}_\Lambda^\Lambda + \mathfrak{a}'$ is closed and the space F_Λ is a core for it. \square

Definition 3.2. *The self-adjoint operator H_ω^Λ in $L^2(\mathbb{R}^2)$ corresponding to the form $\mathfrak{a}_\omega^\Lambda$ via the first representation theorem (see *e.g.* [25, Ch. VI, Thm. 2.1]) is called Schrödinger operator with δ' -interaction of strength ω supported on Λ .*

If ω is a non-negative function, then we occasionally say that the respective δ' -interaction is *attractive*.

3.2. Essential spectrum. In this subsection we characterise the essential spectrum of the operator H_ω^Λ . To this aim we require the following auxiliary lemma.

Lemma 3.3. *Let the self-adjoint operator H_ω^Λ be as in Definition 3.2. Then for any $u \in \mathcal{D}(\mathbb{R}^2 \setminus \Lambda)$ holds*

$$(3.2) \quad u \in \text{dom } H_\omega^\Lambda \quad \text{and} \quad H_\omega^\Lambda u = -\Delta u.$$

Proof. Let Σ and Ω_\pm be as in Hypothesis 2.1. Let $u \in \mathcal{D}(\mathbb{R}^2 \setminus \Lambda) \subset F_\Lambda \subset \text{dom } \mathfrak{a}_\omega^\Lambda$ and $v \in \text{dom } \mathfrak{a}_\omega^\Lambda$. Define $u_\pm := u \upharpoonright \Omega_\pm$ and $v_\pm := v \upharpoonright \Omega_\pm$. With these notations in hands we get

$$\mathfrak{a}_\omega^\Lambda[u, v] = (\nabla u_+, \nabla v_+)_{\Omega_+} + (\nabla u_-, \nabla v_-)_{\Omega_-},$$

where the boundary term in (3.1) vanished due to the choice of u . Applying the first Green identity (see *e.g.* [35, Lem. 4.1] and also [3, Sec. 2]) to the above formula, we get

$$\begin{aligned} \mathfrak{a}_\omega^\Lambda[u, v] &= (-\Delta u_+, v_+)_{\Omega_+} + (-\Delta u_-, v_-)_{\Omega_-} \\ &\quad + (\partial_{\nu_+} u_+|_{\Sigma \setminus \Lambda} + \partial_{\nu_-} u_-|_{\Sigma \setminus \Lambda}, v|_{\Sigma \setminus \Lambda})_{\Sigma \setminus \Lambda} \\ &\quad + (\partial_{\nu_+} u_+|_\Lambda, v_+|_\Lambda)_\Lambda + (\partial_{\nu_-} u_-|_\Lambda, v_-|_\Lambda)_\Lambda = (-\Delta u, v)_{\mathbb{R}^2}, \end{aligned}$$

where we employed that $\partial_{\nu_\pm} u_\pm|_\Lambda = 0$, that $\partial_{\nu_+} u_+|_{\Sigma \setminus \Lambda} + \partial_{\nu_-} u_-|_{\Sigma \setminus \Lambda} = 0$, and that $[v]_{\Sigma \setminus \Lambda} = 0$; for the latter *cf.* Proposition 2.4 (i). Finally, the first representation theorem yields (3.2). \square

Next, we define the notion of the quasi-conical domain; see [22] and also [11, Def. X.6.1].

Definition 3.4. A domain $\Omega \subset \mathbb{R}^2$ is called quasi-conical if for any $n \in \mathbb{N}$ there exists $x_n \in \mathbb{R}^2$ such that $D_n(x_n) \subset \Omega$. Recall that here $D_n(x_n)$ is the disc of radius n with the center x_n .

Using this notion, we prove that positive semi-axis lies inside the spectrum of H_ω^Λ if the domain $\mathbb{R}^2 \setminus \Lambda$ is quasi-conical. The technique of this proof is rather standard.

Proposition 3.5. Let the curve $\Lambda \subset \mathbb{R}^2$ as in Hypotheses 2.1 be such that the domain $\mathbb{R}^2 \setminus \Lambda$ is quasi-conical. Then the spectrum of the self-adjoint operator H_ω^Λ in Definition 3.2 satisfies

$$(3.3) \quad \sigma(H_\omega^\Lambda) \supseteq [0, +\infty).$$

Proof. First, for any $k \in \mathbb{R}^2$, define the sequence

$$u_n(x) := v_n(x) e^{ik \cdot x}, \quad n \in \mathbb{N},$$

where $v_n(x) := n^{-1} v(n^{-1}x)$, $n \in \mathbb{N}$, and v is a non-trivial function in $\mathcal{D}(\mathbb{R}^2)$ with $\text{supp } v \subset D_1$ and such that $\|v\|_{\mathbb{R}^2} = 1$. The prefactor in the definition of v_n is chosen in such a way that also each v_n satisfies $\|v_n\|_{\mathbb{R}^2} = 1$. In fact, we have (by direct computations)

$$(3.4) \quad \|v_n\|_{\mathbb{R}^2} = 1, \quad \|\nabla v_n\|_{\mathbb{R}^2} = \frac{\|\nabla v\|_{\mathbb{R}^2}}{n}, \quad \|\Delta v_n\|_{\mathbb{R}^2} = \frac{\|\Delta v\|_{\mathbb{R}^2}}{n^2}.$$

Secondly, we set

$$w_n(x) := u_n(x - x_n), \quad n \in \mathbb{N},$$

with x_n corresponding to the quasi-conical domain $\mathbb{R}^2 \setminus \Lambda$ according to Definition 3.4. Hence, we get

$$\text{supp } w_n \subset D_n(x_n) \subset \mathbb{R}^2 \setminus \Lambda,$$

and therefore $w_n \in \mathcal{D}(\mathbb{R}^2 \setminus \Lambda)$ for all $n \in \mathbb{N}$. It is clear in view of Lemma 3.3 that each w_n belongs to $\text{dom } H_\omega^\Lambda \supset \mathcal{D}(\mathbb{R}^2 \setminus \Lambda)$.

A direct computation yields

$$|-\Delta w_n - |k|^2 w_n|^2 \leq 2|\Delta v_n|^2 + 4|k \cdot \nabla v_n|^2 \leq 2|\Delta v_n|^2 + 4|k|^2 \cdot |\nabla v_n|^2.$$

Using (3.4) and Lemma 3.3, we therefore have

$$\|H_\omega^\Lambda w_n - |k|^2 w_n\|_{\mathbb{R}^2}^2 = \|-\Delta w_n - |k|^2 w_n\|_{\mathbb{R}^2}^2 \leq 2\|\Delta v_n\|_{\mathbb{R}^2}^2 + 4|k|^2 \|\nabla v_n\|_{\mathbb{R}^2}^2 \rightarrow 0, \quad n \rightarrow \infty.$$

Since the choice of $k \in \mathbb{R}^2$ was arbitrary, we conclude applying Weyl's criterion (see [41, Sec. 7.4] and also [30, Thm. 4]) that $[0, +\infty) \subseteq \sigma(H_\omega^\Lambda)$. \square

We emphasize that not for every non-closed curve $\Lambda \subset \mathbb{R}^2$ the domain $\mathbb{R}^2 \setminus \Lambda$ is quasi-conical; *e.g.* for the Archimedean spiral, defined in polar coordinates (r, φ) by the equation $r(\varphi) := a + b\varphi$, $\varphi \in \mathbb{R}_+$, $a, b > 0$, the domain $\mathbb{R}^2 \setminus \Lambda$ is not of this type.

In the case of bounded curves we show that the essential spectrum of H_ω^Λ coincides with the set $[0, +\infty)$.

Proposition 3.6. *Let the bounded curve $\Lambda \subset \mathbb{R}^2$ be as in Hypothesis 2.1 and let the self-adjoint operator H_ω^Λ be as in Definition 3.2. Then its essential spectrum is characterised as*

$$\sigma_{\text{ess}}(H_\omega^\Lambda) = [0, +\infty).$$

Proof. Let the curve $\Sigma \subset \mathbb{R}^2$ and the domains $\Omega_\pm \subset \mathbb{R}^2$ be as in Hypothesis 2.1, in particular, $\Lambda \subset \Sigma$. Let us also set $c := \|\omega\|_\infty$. Consider the sesquilinear form

$$(3.5) \quad \begin{aligned} \mathfrak{a}_c^\Sigma[u, v] &:= (\nabla u, \nabla v)_{\mathbb{R}^2} - c([u]_\Sigma, [v]_\Sigma)_\Sigma, \\ \text{dom } \mathfrak{a}_c^\Sigma &:= H^1(\Omega_+) \oplus H^1(\Omega_-). \end{aligned}$$

According to [3, Prop. 3.1] the form \mathfrak{a}_c^Σ is closed, densely defined, symmetric, and lower-semibounded in $L^2(\mathbb{R}^2)$. The self-adjoint operator H_c^Σ in $L^2(\mathbb{R}^2)$ representing the form \mathfrak{a}_c^Σ , satisfies

$$(3.6) \quad \sigma_{\text{ess}}(H_c^\Sigma) = [0, +\infty);$$

see [3, Thm. 4.2] and also [5, Thm. 3.16]. The sesquilinear forms $\mathfrak{a}_\omega^\Lambda$ and \mathfrak{a}_c^Σ in (3.1) and (3.5), respectively, naturally satisfy the ordering

$$\mathfrak{a}_c^\Sigma \prec \mathfrak{a}_\omega^\Lambda$$

in the sense of [25, §VI.2.5], see also [6, §10.2.3]. Indeed, firstly, $\text{dom } \mathfrak{a}_\omega^\Lambda \subset \text{dom } \mathfrak{a}_c^\Sigma$ and, secondly, for any $u \in \text{dom } \mathfrak{a}_\omega^\Lambda$ the inequality $\mathfrak{a}_c^\Sigma[u, u] \leq \mathfrak{a}_\omega^\Lambda[u, u]$ holds due to the choice of the constant $c \geq 0$. Hence, using (3.6) and [6, §10.2, Thm. 4] we arrive at

$$0 = \inf \sigma_{\text{ess}}(H_c^\Sigma) \leq \inf \sigma_{\text{ess}}(H_\omega^\Lambda).$$

Therefore, we end up with the following inclusion

$$(3.7) \quad \sigma_{\text{ess}}(H_\omega^\Lambda) \subseteq [0, +\infty).$$

Moreover, for simple geometric reasons for any bounded curve Λ the domain $\mathbb{R}^2 \setminus \Lambda$ is quasi-conical and hence by Proposition 3.5 the opposite inclusion

$$(3.8) \quad \sigma_{\text{ess}}(H_\omega^\Lambda) \supseteq [0, +\infty)$$

holds as well. The claim then follows from these two inclusions ((3.7) and (3.8)). \square

4. NON-NEGATIVITY OF H_ω^Λ

This section plays the central role in the present paper. We obtain various sufficient conditions for the operator H_ω^Λ to be non-negative. Under additional assumptions we also show that positive spectrum of H_ω^Λ comprises the whole positive real axis and thus the operator H_ω^Λ has no bound states. In the proofs we use the min-max principle for self-adjoint operators, a reduction to the one-dimensional model problem discussed in Subsection 2.4, and some insights from geometry and complex analysis.

4.1. An auxiliary lemma. In this subsection we prove a lemma, based on which we show non-negativity of the operators H_ω^Λ under certain assumptions on ω . For the formulation of this lemma we require the following hypothesis, the assumptions of which are grouped in three logical blocks labelled by capital latin letters.

Hypothesis 4.1. (A) Let a monotone piecewise- C^1 curve $\Lambda \subset \mathbb{R}^2$ be parametrized via the mapping $\varphi: (0, R) \rightarrow \mathbb{R}$, $R \in (0, +\infty]$, as in Definition 2.7 with $x_0 = 0$.

(B) Suppose that piecewise- C^1 domains $G_\pm \subset D_R$ satisfy the following conditions:

$$G_+ \cap G_- = \emptyset, \quad \overline{D_R} = \overline{G_+} \cup \overline{G_-}, \quad \text{and} \quad \Lambda \subset \overline{G_+} \cap \overline{G_-}.$$

Set $\Sigma := \overline{G_+} \cap \overline{G_-}$. In particular, the inclusion $\Lambda \subset \Sigma$ holds.

(C) Let the function $\omega \in L^\infty(\Lambda; \mathbb{R})$ as a function of the distance r from the origin satisfy

$$(4.1) \quad \omega(r) \leq \frac{1}{2\pi r \sqrt{1 + (r\varphi'(r))^2}}, \quad \text{for } r \in (0, R).$$

We further deal with the space $H^1(G_+) \oplus H^1(G_-) \subset L^2(D_R)$. Let us introduce also the following notations

$$[u]_\bullet := u_+|_\bullet - u_-|_\bullet, \quad \bullet = \Lambda \text{ or } \Sigma \setminus \Lambda, \quad u = u_+ \oplus u_- \in H^1(G_+) \oplus H^1(G_-).$$

Clearly, one can define polar coordinates (r, φ) on D_R , which are connected with the usual Cartesian coordinates via standard relations $x = r \cos \varphi$ and $y = r \sin \varphi$. The points $(r, \varphi + 2\pi k)$ with $k \in \mathbb{Z}$ are identified with each other. The disc D_R in the polar coordinate system is given by $D_R = \{(r, \varphi) : r \in [0, R), \varphi \in [0, 2\pi)\}$.

For the substantial simplification of further computations we make use of the following short-hand notation:

$$(4.2) \quad \mathfrak{f}_{D_R, \omega}^\Lambda[u] := \|\nabla u\|_{D_R}^2 - (\omega[u]_\Lambda, [u]_\Lambda)_\Lambda, \quad u \in \mathcal{D}(\overline{G_+}) \oplus \mathcal{D}(\overline{G_-}),$$

where all the objects are as in Hypothesis 4.1. Now we formulate and prove the following lemma.

Lemma 4.1. *Assume that Hypothesis 4.1 holds. Then $\mathfrak{f}_{D_R, \omega}^\Lambda[u] \geq 0$ for all $u \in \mathcal{D}(\overline{G_+}) \oplus \mathcal{D}(\overline{G_-})$ such that $[u]_{\Sigma \setminus \Lambda} = 0$.*

Proof. Let $u \in \mathcal{D}(\overline{G_+}) \oplus \mathcal{D}(\overline{G_-})$ be such that $[u]_{\Sigma \setminus \Lambda} = 0$. The proof of $\mathfrak{f}_{D_R, \omega}^\Lambda[u] \geq 0$ is then split in three steps.

Step 1. For any $(x, y) \in D_R \setminus \Sigma$ the value $|(\nabla u)(x, y)|^2$ can be expressed in polar coordinates (r, φ) as

$$|(\nabla u)(x, y)|^2 = |(\partial_r u)(r, \varphi)|^2 + \frac{1}{r^2} |(\partial_\varphi u)(r, \varphi)|^2.$$

Using the above expression for the gradient we obtain the following estimate

$$(4.3) \quad \|\nabla u\|_{D_R}^2 = \int_0^{2\pi} \int_0^R |(\nabla u)(r, \varphi)|^2 r dr d\varphi \leq \int_0^R \frac{1}{r} \left(\int_0^{2\pi} |(\partial_\varphi u)(r, \varphi)|^2 d\varphi \right) dr,$$

in which we have thrown away a positive term in the second step. Interchanging of the integrals in the above computation can be justified by Fubini's theorem (see *e.g.* [40, Ch. 2, Thm. 3.1]).

Step 2. Using the mapping $\varphi: (0, R) \rightarrow \mathbb{R}$ as in Hypothesis 4.1 (A) we define the following auxiliary function

$$\mathfrak{j}(r) := \sqrt{1 + (r\varphi'(r))^2}, \quad r \in (0, R).$$

The curvilinear integral along Λ in (4.2) can be rewritten in terms of the mapping $\varphi: (0, R) \rightarrow \mathbb{R}$ and the function \mathfrak{j} in the conventional way and then further estimated with the help of assumption (4.1)

$$(4.4) \quad \begin{aligned} (\omega[u]_\Lambda, [u]_\Lambda)_\Lambda &= \int_0^R \omega(r) \mathfrak{j}(r) |u_+(r, \varphi(r)) - u_-(r, \varphi(r))|^2 dr \\ &\leq \int_0^R \frac{1}{2\pi r} |u_+(r, \varphi(r)) - u_-(r, \varphi(r))|^2 dr. \end{aligned}$$

Step 3. Define the following function

$$(4.5) \quad S(r) := \int_0^{2\pi} |(\partial_\varphi u)(r, \varphi)|^2 d\varphi - \frac{1}{2\pi} |u_+(r, \varphi(r)) - u_-(r, \varphi(r))|^2,$$

where $r \in (0, R)$. Thanks to the choice of u , for all $r \in (0, R)$ the function $[0, 2\pi) \ni \varphi \mapsto u(r, \varphi)$ can naturally be identified with the piecewise- C^1 function ψ_r on the interval $I = (0, 2\pi)$, which by [36, App. E] belongs to $H^1(I)$. Moreover, the relation $S(r) = \mathfrak{a}_{d,\omega}[\psi_r]$ holds with the form $\mathfrak{a}_{d,\omega}$ as in (2.11), where $d = 2\pi$ and $\omega = 1/2\pi$. In particular, $d\omega = 2\pi/2\pi = 1$ and by Corollary 2.12 we obtain

$$S(r) \geq 0, \quad \text{for all } r \in (0, R).$$

Finally, using (4.3), (4.4) and non-negativity of $S(r)$ we arrive at

$$\mathfrak{f}_{D_R,\omega}^\Lambda[u] \geq \int_0^R \frac{S(r)}{r} dr \geq 0.$$

□

4.2. Non-negativity of H_ω^Λ for monotone Λ . In this subsection we obtain various explicit sufficient conditions on ω ensuring non-negativity of H_ω^Λ assuming that Λ is monotone. General results are illustrated with two examples: an Archimedean spiral and a subinterval of the straight line in \mathbb{R}^2 .

Theorem 4.2. *Let a monotone piecewise- C^1 curve $\Lambda \subset \mathbb{R}^2$ be parametrized via $\varphi: (0, R) \rightarrow \mathbb{R}$, $R \in (0, +\infty]$, as in Definition 2.7. Let the self-adjoint operator H_ω^Λ be as in Definition 3.2 with $\omega \in L^\infty(\Lambda; \mathbb{R})$. Then*

$$\sigma(H_\omega^\Lambda) \subseteq [0, +\infty) \quad \text{if } \omega(r) \leq \frac{1}{2\pi r \sqrt{1 + (r\varphi'(r))^2}}, \quad \text{for } r \in (0, R).$$

If ω is majorized as above, and additionally, the domain $\mathbb{R}^2 \setminus \Lambda$ is quasi-conical, then $\sigma(H_\omega^\Lambda) = [0, +\infty)$.

Proof. Let Σ and Ω_\pm be as in Hypothesis 2.1. Without loss of generality we assume that $\mathbf{x}_0 = 0$ in Definition 2.7.

Let us define the complement $\Omega_c := \mathbb{R}^2 \setminus \overline{D_R}$ of the disc D_R , the curve $\Gamma := \Sigma \cap D_R$, and the domains $G_\pm := \Omega_\pm \cap D_R$. It is straightforward to see that the tuple $\{D_R, G_+, G_-, \Lambda, \omega\}$ satisfies Hypothesis 4.1.

Let $u \in F_\Lambda$ and define $u_R := u \upharpoonright D_R$, $u_c := u \upharpoonright \Omega_c$. Then it holds that

$$u_R \in \mathcal{D}(\overline{G_+}) \oplus \mathcal{D}(\overline{G_-}) \quad \text{and} \quad [u_R]_{\Gamma \setminus \Lambda} = 0.$$

Hence, using Lemma 4.1 we get

$$\mathfrak{a}_\omega^\Lambda[u, u] = \mathfrak{f}_{D_R,\omega}^\Lambda[u_R] + \|\nabla u_c\|_{\Omega_c}^2 \geq \mathfrak{f}_{D_R,\omega}^\Lambda[u_R] \geq 0.$$

Since F_Λ is a core for the form $\mathfrak{a}_\omega^\Lambda$, we get by [8, Thm. 4.5.3] that the self-adjoint operator H_ω^Λ is non-negative. If, additionally, the domain $\mathbb{R}^2 \setminus \Lambda$ is quasi-conical, Proposition 3.6 implies that

$$\sigma(H_\omega^\Lambda) = [0, +\infty).$$

□

Example 4.3. *Let the piecewise- C^1 curve $\Lambda \subset \mathbb{R}^2$ be defined as*

$$\Lambda := \{(r \cos(r), r \sin(r)) \in \mathbb{R}^2 : r \in \mathbb{R}_+\}.$$

Obviously, this curve is monotone in the sense of Definition 2.7 with $\mathbf{x}_0 = 0$ and $\varphi(r) := r$, $r \in (0, +\infty)$. The curve Λ is a special case of an Archimedean spiral. Theorem 4.2 yields that

$$\sigma(H_\omega^\Lambda) \subseteq [0, +\infty) \quad \text{if } \omega(r) \leq \frac{1}{2\pi r \sqrt{1 + r^2}}, \quad \text{for } r > 0.$$

The case of a non-varying interaction strength ω is of special interest. In the rest of this subsection we assume for the sake of demonstrativeness that $\omega \in \mathbb{R}$ is a constant. Define also the following characteristic of a bounded monotone piecewise- C^1 curve $\Lambda \subset \mathbb{R}^2$ (parametrized as in Definition 2.7)

$$(4.6) \quad \omega_*(\Lambda) := \inf_{r \in (0, R)} \frac{1}{2\pi r \sqrt{1 + (r\varphi'(r))^2}}.$$

It is not difficult to see that $0 < \omega_*(\Lambda) < +\infty$.

The following corollary is a direct consequence of Theorem 4.2, Proposition 3.6, and simple geometric argumentation.

Corollary 4.4. *Let $\Lambda \subset \mathbb{R}^2$ be a bounded monotone piecewise- C^1 curve and let the self-adjoint operator H_ω^Λ be as in Definition 3.2 with non-varying strength $\omega \in \mathbb{R}$. Then*

$$\sigma(H_\omega^\Lambda) = [0, +\infty) \quad \text{for all } \omega \in (-\infty, \omega_*]$$

with $\omega_* = \omega_*(\Lambda) > 0$ defined in (4.6).

To illustrate this corollary we provide an example.

Example 4.5. *Consider the interval of length $L > 0$ in the plane:*

$$(4.7) \quad \Lambda := \{(x, 0) \in \mathbb{R}^2 : 0 < x < L\}.$$

Clearly, the interval Λ is monotone in the sense of Definition 2.7 with $x_0 = 0$ and $\varphi(r) = 0$, $r \in (0, L)$. Then we get from Corollary 4.4, using formula (4.6), that

$$\sigma(H_\omega^\Lambda) = [0, +\infty) \quad \text{for all } \omega \in (-\infty, \frac{1}{2\pi L}].$$

Remark 4.6. Let Λ be as in (4.7). It is worth noting that the result of the above example can be improved in the following way. Define the points $x_0 = (0, 0)$, $x_1 = (L, 0)$, the intervals

$$\Lambda_0 := \{(x, 0) \in \mathbb{R}^2 : 0 < x < L/2\}, \quad \Lambda_1 := \{(x, 0) \in \mathbb{R}^2 : L/2 < x < L\},$$

the discs $D_{L/2}(x_0)$, $D_{L/2}(x_1)$, and the complement

$$\Omega_c := \mathbb{R}^2 \setminus (\overline{D_{L/2}(x_0)} \cup \overline{D_{L/2}(x_1)})$$

of the closure of their union. Let $u \in F_\Lambda$ and define $u_k := u \upharpoonright D_{L/2}(x_k)$, $k = 0, 1$, $u_c := u \upharpoonright \Omega_c$. Assuming that $\omega \in (-\infty, \frac{1}{\pi L}]$, we get by Lemma 4.1 that

$$\mathfrak{a}_\omega^\Lambda[u, u] = \mathfrak{f}_{D_{L/2}(x_0), \omega}^{\Lambda_0}[u_0] + \mathfrak{f}_{D_{L/2}(x_1), \omega}^{\Lambda_1}[u_1] + \|\nabla u_c\|_{\Omega_c}^2 \geq 0.$$

Thus, the operator H_ω^Λ is non-negative and by Proposition 3.6 we get $\sigma(H_\omega^\Lambda) = [0, +\infty)$.

One may expect that for a sufficiently large coupling constant $\omega > 0$ or for a sufficiently long curve Λ negative spectrum of the self-adjoint operator H_ω^Λ is non-empty. In the next proposition we confirm this expectation via an example.

Proposition 4.7. *Let $\Lambda \subset \mathbb{R}^2$ be as in (4.7) and the self-adjoint operator H_ω^Λ be as in Definition 3.2 with non-varying strength $\omega \in \mathbb{R}$. Then*

$$\sigma_d(H_\omega^\Lambda) \cap (-\infty, 0) \neq \emptyset \quad \text{for all } \omega \in (\frac{\pi}{2L}, +\infty).$$

Proof. Let us split the plane \mathbb{R}^2 into three domains

$$\Omega_1 := (-\infty, 0) \times \mathbb{R}, \quad \Omega_2 := (0, L) \times \mathbb{R}, \quad \Omega_3 := (L, +\infty) \times \mathbb{R},$$

via straight lines

$$\Pi_1 := \{0\} \times \mathbb{R}, \quad \Pi_2 := \{L\} \times \mathbb{R},$$

as indicated in Figure 4.1.

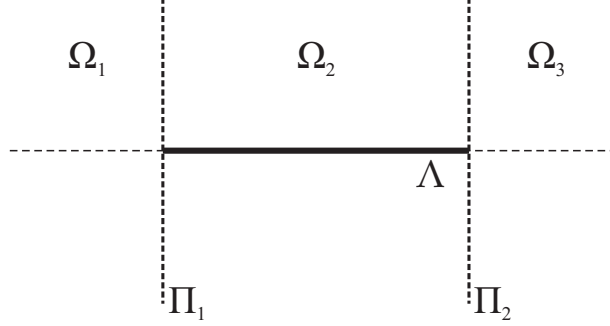


FIGURE 4.1. Splitting of \mathbb{R}^2 into three domains $\{\Omega_k\}_{k=1}^3$.

Consider the sesquilinear form

$$\mathfrak{a}_{\omega, D}^\Lambda[u, v] := \mathfrak{a}_\omega^\Lambda[u, v], \quad \text{dom } \mathfrak{a}_{\omega, D}^\Lambda := \{u \in \text{dom } \mathfrak{a}_\omega^\Lambda : u|_{\Pi_k} = 0, \ k = 1, 2\}.$$

It is not difficult to check that the sesquilinear form $\mathfrak{a}_{\omega, D}^\Lambda$ is closed, symmetric, densely defined, and semibounded in $L^2(\mathbb{R}^2)$. This form induces via the first representation theorem the self-adjoint operator $H_{\omega, D}^\Lambda$ in $L^2(\mathbb{R}^2)$, which can be represented as the orthogonal sum $H_1 \oplus H_2 \oplus H_3$ with respect to the decomposition $L^2(\mathbb{R}^2) = \oplus_{k=1}^3 L^2(\Omega_k)$. Note that H_1 and H_3 are both non-negative and their spectra are given by the set $[0, +\infty)$. The spectrum of H_2 can be computed via separation of variables in the strip Ω_2 . In particular, the ground state of H_2 corresponds to the eigenvalue

$$\lambda_1(H_2) = \frac{\pi^2}{L^2} - 4\omega^2,$$

where we used that the one-dimensional Schrödinger operator on the full-line with one-center point δ' -interaction of strength $\omega > 0$ has the lowest eigenvalue $-4\omega^2$; cf. [1, Ch. I.4], where not ω , but $\beta = 1/\omega$ is called the strength of δ' -interaction.

If the assumption in the formulation of the proposition holds, then $\lambda_1(H_2) < 0$ and the operator $H_{\omega, D}^\Lambda$ has at least one negative eigenvalue.

It remains to note that by Proposition 3.6 we have $\sigma_{\text{ess}}(H_\omega^\Lambda) = [0, +\infty)$ and that the form ordering

$$\mathfrak{a}_\omega^\Lambda \prec \mathfrak{a}_{\omega, D}^\Lambda$$

can easily be verified, which yields by [6, §10.2, Thm. 4] that the operator H_ω^Λ has at least one negative eigenvalue. \square

4.3. Absence of bound states for H_ω^Λ and LFTs. In this subsection we show absence of bound states in the weak coupling regime for a class of bounded non-monotone piecewise- C^1 curves, which are (with minor restrictions) images of bounded monotone curves under LFTs. Since the identical transform $M(z) = z$ is an LFT, this class is certainly larger than the class of bounded monotone curves. As an example we treat δ' -interaction supported on a circular arc subtending an angle $\theta > \pi$.

First, we provide for convenience of the reader two standard claims on change of variables under LFT. The proofs of them are outsourced to Appendix A.

Lemma 4.8. *Let $\Lambda \subset \mathbb{R}^2$ be a bounded curve as in Hypothesis 2.1, let the space F_Λ be as in (2.1), and let $M: \widehat{\mathbb{C}} \rightarrow \widehat{\mathbb{C}}$ be an LFT such that $M(\infty), M^{-1}(\infty) \notin \Lambda$. Then for any $u \in F_\Lambda$*

$$(4.8) \quad \int_{\mathbb{R}^2} |\nabla u|^2 dx = \int_{\mathbb{R}^2} |\nabla v|^2 dx$$

holds with $v := u \circ M$.

Remark 4.9. The function v itself in the formulation of the above lemma is continuous and piecewise smooth, but it is not necessarily compactly supported or square-integrable.

Lemma 4.10. *Let $\Lambda \subset \mathbb{R}^2$ be a bounded piecewise- C^1 curve, parametrized via the mapping $\lambda: I \rightarrow \mathbb{R}^2$, $I := (0, L)$, as in Definition 2.6, let the space F_Λ be as in (2.1) and let $\omega \in L^\infty(\Lambda; \mathbb{R})$. For an LFT $M: \widehat{\mathbb{C}} \rightarrow \widehat{\mathbb{C}}$ with J_M as in (2.6) and such that $M(\infty), M^{-1}(\infty) \notin \Lambda$ define $\Gamma := M^{-1}(\Lambda)$, $\gamma := M^{-1} \circ \lambda$ and*

$$(4.9) \quad \tilde{\omega}(\gamma(s)) := \omega(\lambda(s)) \sqrt{J_M(\gamma(s))}, \quad s \in I.$$

Then the relation

$$(\omega[u]_\Lambda, [u]_\Lambda)_\Lambda = (\tilde{\omega}[v]_\Gamma, [v]_\Gamma)_\Gamma$$

holds for any $u \in F_\Lambda$ and $v := u \circ M$.

Remark 4.11. Note that the function v in the formulation of the above lemma does not belong to F_Γ in general. However, $v_\pm := v \upharpoonright M^{-1}(\Omega_\pm)$ with Ω_\pm as in Hypothesis 2.1 are well-defined and continuous up to Γ . Hence, the restrictions $v_\pm|_\Gamma$ are meaningful and $[v]_\Gamma := v_+|_\Gamma - v_-|_\Gamma$ is well-defined.

Now we can formulate the key result of this subsection, whose proof with all the above preparations is rather short.

Theorem 4.12. *Let $\Lambda \subset \mathbb{R}^2$ be a bounded piecewise- C^1 curve and let the self-adjoint operator H_ω^Λ in $L^2(\mathbb{R}^2)$ be as in Definition 3.2 with non-varying strength $\omega \in \mathbb{R}$. Suppose that there exists an LFT $M: \widehat{\mathbb{C}} \rightarrow \widehat{\mathbb{C}}$ such that:*

- (a) $M(\infty), M^{-1}(\infty) \notin \Lambda$;
- (b) $\Gamma := M^{-1}(\Lambda)$ is monotone.

Let the constant $\omega_(\Gamma) > 0$ be associated to Γ via (4.6). Then it holds that*

$$\sigma(H_\omega^\Lambda) = [0, +\infty) \quad \text{for all } \omega \in (-\infty, \omega_*],$$

where $\omega_ := \frac{\omega_*(\Gamma)}{\sup_{z \in \Gamma} \sqrt{J_M(z)}}$.*

Proof. Step 1. Suppose that the curve Λ is parametrized via the mapping $\lambda: (0, L) \rightarrow \mathbb{R}^2$ (as in Definition 2.6). Define the mapping $\gamma := M^{-1} \circ \lambda$. Due to assumption (a) the curve Γ is bounded and the mapping γ parametrizes it. Without loss of generality suppose that the curve Γ

is monotone in the sense of Definition 2.7 with $x_0 = \gamma(0) = 0$ and with $\varphi: (0, R) \rightarrow \mathbb{R}$, $R = |\gamma(L)|$. Consider the complement $\Omega_c := \mathbb{R}^2 \setminus \overline{D_R}$ of the disc D_R . Let the curve Σ and the domains Ω_\pm be associated to Λ as in Hypothesis 2.1.

Define auxiliary domains $G_\pm := M^{-1}(\Omega_\pm) \cap D_R$. Thus, the splitting

$$D_R = G_+ \dot{\cup} M^{-1}(\Sigma) \dot{\cup} G_-$$

holds. Let $\tilde{\omega}$ be defined via the formula (4.9) in Lemma 4.10. Hence, we obtain

$$\tilde{\omega} \leq \omega \sup_{z \in \Gamma} \sqrt{J_M(z)} \leq \omega_* \sup_{z \in \Gamma} \sqrt{J_M(z)} = \omega_*(\Gamma).$$

Summarizing, the tuple $\{D_R, G_+, G_-, \Gamma, \tilde{\omega}\}$ fulfils Hypothesis 4.1.

Step 2. Let $u \in F_\Lambda$ with F_Λ as in (2.1) and define the composition $v := u \circ M$. Set $v_R := v \upharpoonright D_R$ and $v_c := v \upharpoonright \Omega_c$. Using Lemmas 4.8 and 4.10 we obtain

$$(4.10) \quad \begin{aligned} \mathfrak{a}_\omega^\Lambda[u, u] &= \|\nabla u\|_{\mathbb{R}^2}^2 - (\omega[u]_\Lambda, [u]_\Lambda)_\Lambda = \|\nabla v\|_{\mathbb{R}^2}^2 - (\tilde{\omega}[v]_\Gamma, [v]_\Gamma)_\Gamma \\ &= \|\nabla v_R\|_{D_R}^2 - (\tilde{\omega}[v_R]_\Gamma, [v_R]_\Gamma)_\Gamma + \|\nabla v_c\|_{\Omega_c}^2 \geq \mathfrak{f}_{D_R, \tilde{\omega}}^\Gamma[v_R, v_R] \geq 0, \end{aligned}$$

where we applied Lemma 4.1 in the last step. Hence, the operator H_ω^Λ is non-negative.

Step 3. Since the curve Λ is bounded, Proposition 3.6 applies, and we arrive at $\sigma_{\text{ess}}(H_\omega^\Lambda) = [0, +\infty)$. The results of Step 2 and Step 3 imply the claim. \square

To conclude this subsection we show that a model of sufficiently weak δ' -interaction of non-varying strength supported on a circular arc subtending the angle $2\pi - 2\varepsilon$ ($\varepsilon \in (0, \pi)$) has no bound states in the weak coupling regime. We emphasize that circular arcs subtending angles $\theta > \pi$ are non-monotone and the results of the previous subsection do not apply to them.

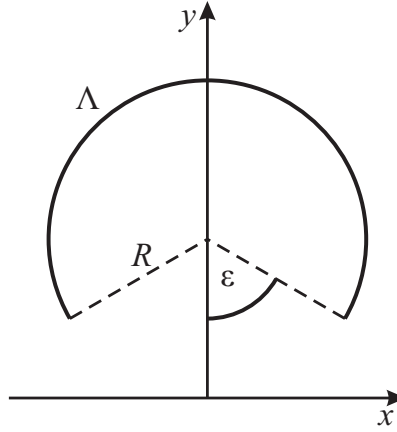


FIGURE 4.2. The circular arc of radius $R > 0$ subtending the angle $2\pi - 2\varepsilon$ with $\varepsilon \in (0, \pi)$.

Example 4.13. The circular arc (see Figure 4.2) can be efficiently parametrized as follows:

$$(4.11) \quad \Lambda := \{(R \sin \varphi, R(1 - \cos \varphi)) \in \mathbb{R}^2 : \varphi \in (\varepsilon, 2\pi - \varepsilon)\},$$

where $\varepsilon \in (0, \pi)$ and $R > 0$ is the radius of the underlying circle. Consider the LFT $M(z) := 1/z$. One easily sees that

$$M_1(x, y) = \operatorname{Re} M(x + iy) = \frac{x}{x^2 + y^2},$$

and according to (2.6) the Jacobian J_M of this LFT is given by the formula

$$(4.12) \quad \begin{aligned} J_M(x, y) &= ((\partial_x M_1)^2 + (\partial_y M_1)^2)(x, y) \\ &= \frac{(x^2 - y^2)^2}{(x^2 + y^2)^4} + \frac{4x^2 y^2}{(x^2 + y^2)^4} = \frac{1}{(x^2 + y^2)^2}. \end{aligned}$$

Next observe that $M(\infty) = M^{-1}(\infty) = 0 \notin \Lambda$. Moreover, this LFT is inverse to itself and under the LFT $M^{-1}(z) = 1/z$ the arc $\Lambda \subset \mathbb{R}^2$ is mapped onto the interval

$$\Gamma := M^{-1}(\Lambda) = \{(x, -\frac{1}{2R}) \in \mathbb{R}^2 : |x| < \cot(\varepsilon/2)/(2R)\},$$

which is obviously monotone in the sense of Definition 2.7. Compute further $\omega_*(\Gamma)$ defined in (4.6)

$$\omega_*(\Gamma) = \inf_{r \in (0, |\Gamma|)} \frac{1}{2\pi r} = \frac{1}{2\pi |\Gamma|} = \frac{R}{2\pi \cot(\varepsilon/2)};$$

here $|\Gamma|$ is the length of Γ . Moreover, we obtain from (4.12) that

$$\sup_{z \in \Gamma} \sqrt{J_M(z)} = 4R^2.$$

Hence, Theorem 4.12 implies that

$$\sigma(H_\omega^\Lambda) = [0, +\infty) \quad \text{for all } \omega \in (-\infty, (8\pi R)^{-1} \tan(\varepsilon/2)].$$

5. REMARKS AND OPEN QUESTIONS

In the present paper we have analysed from various perspectives a new effect of absence of the negative spectrum for Hamiltonians with δ' -interaction supported on non-closed curves in \mathbb{R}^2 . Quite a few questions remain open and we wish to formulate two of them.

Comparing Example 4.5 and Proposition 4.7 one may pose the following question.

Open Question A. *Let the constant $L > 0$ be fixed and the interval Λ be as in (4.7). The problem is to find the critical strength $\omega_{\text{cr}}(L) > 0$ such that the operator H_ω^Λ is non-negative if and only if $\omega \in \mathbb{R}$ satisfies $\omega \leq \omega_{\text{cr}}(L)$.*

The same question as above can be asked for other shapes of Λ , but the authors do not expect that an exact formula for the critical strength can be found.

On one hand, our method of the proof does not allow to cover curves of generic shape. On the other hand, despite many attempts, we have not found out any example of a bounded non-closed curve, for which bound states in the weak coupling regime do exist. A general open question can be posed.

Open Question B. *Is it true that for any bounded sufficiently smooth non-closed curve $\Lambda \subset \mathbb{R}^2$ there exists a constant $\omega_* > 0$ such that $\sigma(H_\omega^\Lambda) = [0, +\infty)$ for all $\omega \in (-\infty, \omega_*]$?*

It is worth noting that the program carried out in Subsection 4.3 for linear fractional transformations can be generalized by means of Neumann bracketting to arbitrary conformal maps. This could be a possible way to answer Question B.

Finally, we mention that several assumptions play only technical role and can be removed with additional efforts. Namely, assuming that Λ is a subarc of the boundary of a Lipschitz domain is technical as well assuming that the curve Λ is piecewise- C^1 in some of the formulations instead of just being Lipschitz.

APPENDIX A. PROOFS OF LEMMAS 4.8 AND 4.10

Proof of Lemma 4.8. Consider the following two open sets

$$E := \mathbb{R}^2 \setminus (\{M(\infty)\} \cup M(\Lambda)) \quad \text{and} \quad F := \mathbb{R}^2 \setminus \Lambda.$$

By the formula in Lemma 2.10 and using that $\mathbb{R}^2 \setminus E$ is a null set we get

$$\int_{\mathbb{R}^2} |\nabla v|^2 dx = \int_E |\nabla v|^2 dx = \int_E |(\nabla u) \circ M|^2 J_M dx.$$

According to Proposition 2.9 we have that $M^{-1}: E \rightarrow F$ is a bijection which is additionally everywhere differentiable in E ; cf. (2.5). Hence, we can apply the substitution rule for Lebesgue integrals (e.g. [31, Thm. 8.21, Cor. 8.22]) and get

$$\begin{aligned} \int_{\mathbb{R}^2} |\nabla v|^2 dx &= \int_F |(\nabla u) \circ M \circ M^{-1}|^2 J_M (J_M)^{-1} dx \\ &= \int_F |\nabla u|^2 dx = \int_{\mathbb{R}^2} |\nabla u|^2 dx; \end{aligned}$$

where in the last step we employed that $\mathbb{R}^2 \setminus F$ is a null set. \square

Proof of Lemma 4.10. Observe first that by definition of the curvilinear integral we have

$$(A.1) \quad (\tilde{\omega}[v]_\Gamma, [v]_\Gamma)_\Gamma = \int_0^L \tilde{\omega}(\gamma(s)) |v_+(\gamma(s)) - v_-(\gamma(s))|^2 |\gamma'(s)| ds.$$

Using elementary composition rules we also note

$$(A.2) \quad \begin{aligned} v_+(\gamma(s)) - v_-(\gamma(s)) &= (u_+ \circ M \circ M^{-1} \circ \lambda)(s) - (u_- \circ M \circ M^{-1} \circ \lambda)(s) \\ &= u_+(\lambda(s)) - u_-(\lambda(s)), \end{aligned}$$

where $u_\pm = u \upharpoonright \Omega_\pm$ and $v_\pm = v \upharpoonright M^{-1}(\Omega_\pm)$. Observe also that $\lambda = M \circ \gamma$. Using (2.6) and (2.7) we obtain

$$\begin{aligned} |\lambda'(s)|^2 &= ((\nabla M_1 \circ \gamma)(s) \cdot \gamma'(s))^2 + ((\nabla M_2 \circ \gamma)(s) \cdot \gamma'(s))^2 \\ &= |(\nabla M_1 \circ \gamma)(s)|^2 \cdot |\gamma'(s)|^2 \cos^2 \alpha + |(\nabla M_2 \circ \gamma)(s)|^2 \cdot |\gamma'(s)|^2 \sin^2 \alpha \\ &= J_M(\gamma(s)) \cdot |\gamma'(s)|^2, \end{aligned}$$

where α is the angle between ∇M_1 and γ' . Thanks to (A.1) and using (A.2) we arrive at

$$\begin{aligned} (\tilde{\omega}[v]_\Gamma, [v]_\Gamma)_\Gamma &= \int_0^L \tilde{\omega}(\gamma(s)) |v_+(\gamma(s)) - v_-(\gamma(s))|^2 |\gamma'(s)| ds \\ &= \int_0^L \frac{\tilde{\omega}(\gamma(s))}{\sqrt{J_M(\gamma(s))}} |u_+(\lambda(s)) - u_-(\lambda(s))|^2 |\lambda'(s)| ds. \end{aligned}$$

Finally, employing (4.9) we end up with the desired relation

$$(\tilde{\omega}[v]_\Gamma, [v]_\Gamma)_\Gamma = (\omega[u]_\Lambda, [u]_\Lambda)_\Lambda.$$

\square

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